

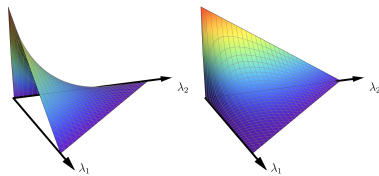
Blow-up Finite Elements

Yakov Berchenko-Kogan, joint with Evan Gawlik

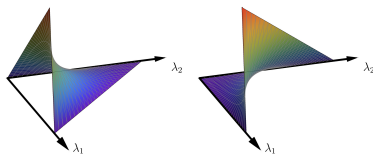
Florida Institute of Technology

April 19–20, 2024

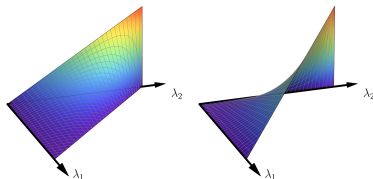
New finite element space



$$\psi_{012} = \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2}, \quad \psi_{021} = \frac{\lambda_0 \lambda_2}{\lambda_2 + \lambda_1},$$

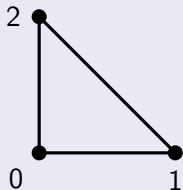


$$\psi_{102} = \frac{\lambda_1 \lambda_0}{\lambda_0 + \lambda_2}, \quad \psi_{120} = \frac{\lambda_1 \lambda_2}{\lambda_2 + \lambda_0},$$



$$\psi_{201} = \frac{\lambda_2 \lambda_0}{\lambda_0 + \lambda_1}, \quad \psi_{210} = \frac{\lambda_2 \lambda_1}{\lambda_1 + \lambda_0}.$$

Classical \mathcal{P}_1

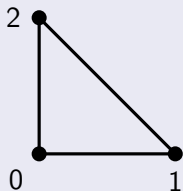


Barycentric coordinates: $\lambda_0 + \lambda_1 + \lambda_2 = 1$.

- 0 : $\lambda_0 = 1 \Leftrightarrow \lambda_1 = \lambda_2 = 0$
- 1 : $\lambda_1 = 1 \Leftrightarrow \lambda_2 = \lambda_0 = 0$
- 2 : $\lambda_2 = 1 \Leftrightarrow \lambda_0 = \lambda_1 = 0$

Degrees of freedom

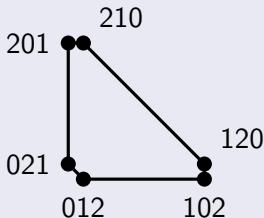
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Blow-up $b\mathcal{P}_1$



- 012 : $\lim_{\lambda_1 \rightarrow 0} \lim_{\lambda_2 \rightarrow 0}$
- 120 : $\lim_{\lambda_2 \rightarrow 0} \lim_{\lambda_0 \rightarrow 0}$
- 201 : $\lim_{\lambda_0 \rightarrow 0} \lim_{\lambda_1 \rightarrow 0}$
- 021 : $\lim_{\lambda_2 \rightarrow 0} \lim_{\lambda_1 \rightarrow 0}$
- 102 : $\lim_{\lambda_0 \rightarrow 0} \lim_{\lambda_2 \rightarrow 0}$
- 210 : $\lim_{\lambda_1 \rightarrow 0} \lim_{\lambda_0 \rightarrow 0}$

Example: Evaluating degrees of freedom

Recall

$$\lambda_0 + \lambda_1 + \lambda_2 = 1, \quad \psi_{012} = \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2}.$$

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$$\lambda_0 + \lambda_1 + \lambda_2 = 1, \quad \psi_{012} = \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2}.$$

Evaluating degrees of freedom

$$012 : \lim_{\lambda_1 \rightarrow 0} \lim_{\lambda_2 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2} = \lim_{\lambda_1 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1} = \lim_{\lambda_0 \rightarrow 1} \lambda_0 = 1,$$

$$021 : \lim_{\lambda_2 \rightarrow 0} \lim_{\lambda_1 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2} = \lim_{\lambda_2 \rightarrow 0} \frac{0}{\lambda_2} = 0,$$

$$120 : \lim_{\lambda_2 \rightarrow 0} \lim_{\lambda_0 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2} = \lim_{\lambda_2 \rightarrow 0} \frac{0}{1} = 0,$$

$$102 : \lim_{\lambda_0 \rightarrow 0} \lim_{\lambda_2 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2} = \lim_{\lambda_0 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1} = 0,$$

$$201 : \lim_{\lambda_0 \rightarrow 0} \lim_{\lambda_1 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2} = \lim_{\lambda_0 \rightarrow 0} \frac{0}{\lambda_2} = 0,$$

$$210 : \lim_{\lambda_1 \rightarrow 0} \lim_{\lambda_0 \rightarrow 0} \frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2} = \lim_{\lambda_1 \rightarrow 0} \frac{0}{1} = 0.$$

Motivating problem

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- Goal: construct **intrinsic** discretizations of tangent vector fields on smooth surfaces that are **continuous across edges**.

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Remark about FEEC

- FEEC discretizations are **intrinsic** but only tangentially continuous across edges. **Normal components** are generally **discontinuous**.

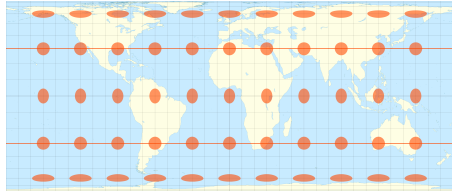
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Remark about FEEC

- FEEC discretizations are **intrinsic** but only tangentially continuous across edges. **Normal components** are generally **discontinuous**.
- FEEC discretization suffices for Hodge Laplacian, but not for Bochner Laplacian.

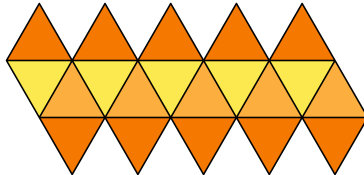
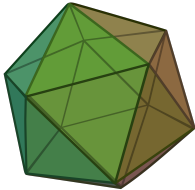
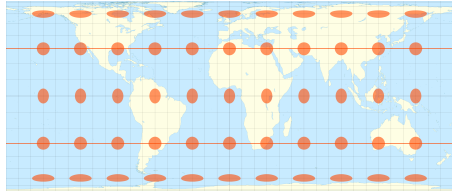
Extrinsic vs. Intrinsic



Four images from Wikipedia

Yakov Berchenko-Kogan, joint with Evan Gawlik

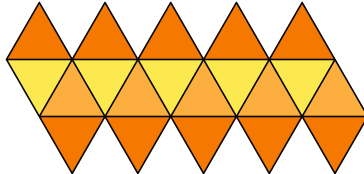
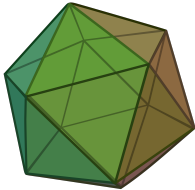
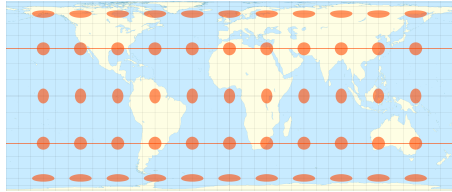
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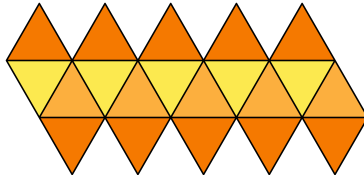
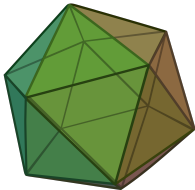
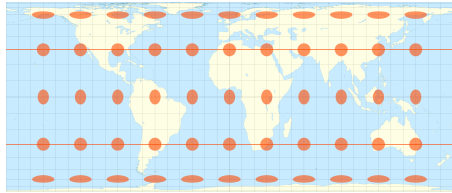


Why compute intrinsically?

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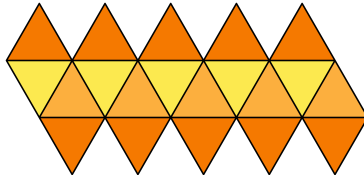
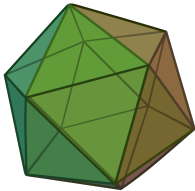
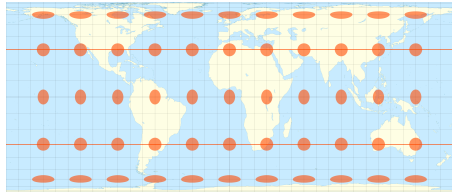


Why compute intrinsically?

- Intrinsic problems, e.g. numerical relativity, Ricci flow.

Four images from Wikipedia

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Why compute intrinsically?

- Intrinsic problems, e.g. numerical relativity, Ricci flow.
- Structure preservation: independence of embedding.

Four images from Wikipedia

Angle defect obstruction to continuous elements

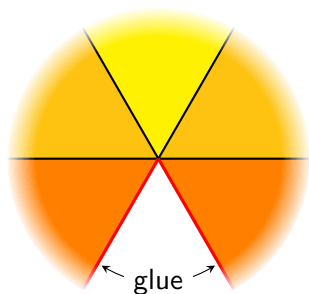
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- What do we see when we zoom in on a vertex?

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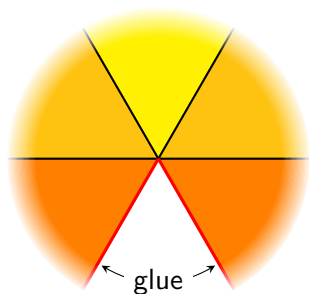
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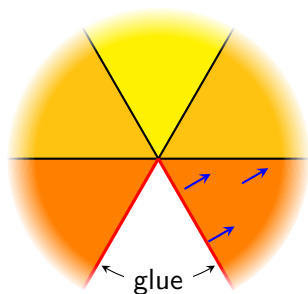
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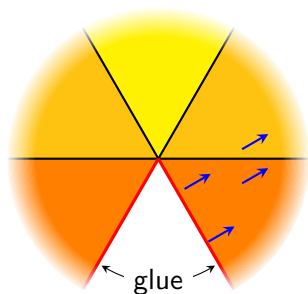
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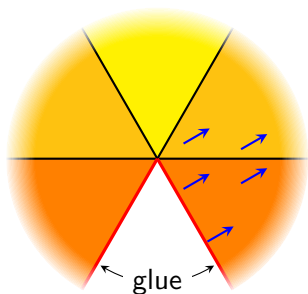
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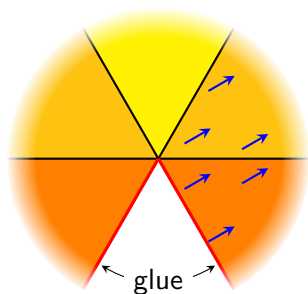
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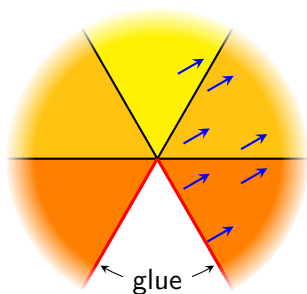
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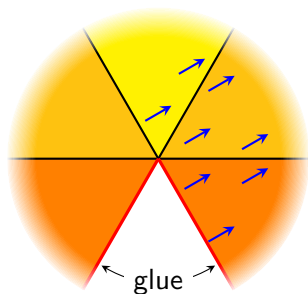
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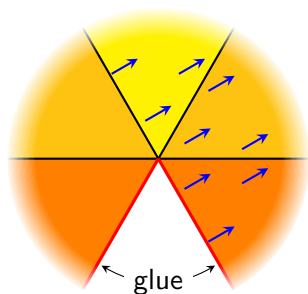
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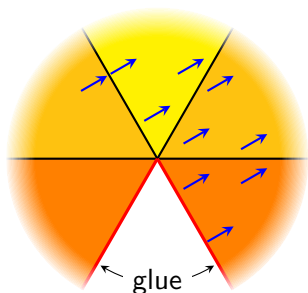
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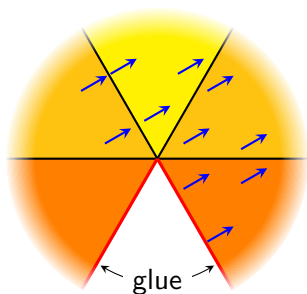
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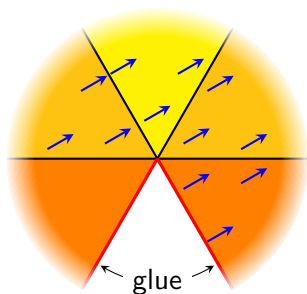
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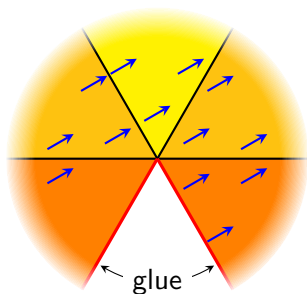
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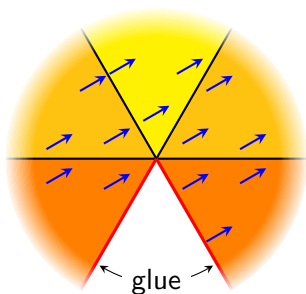
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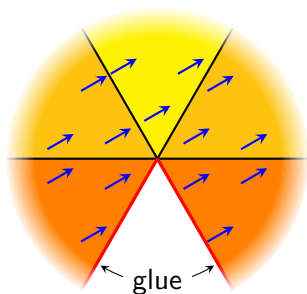
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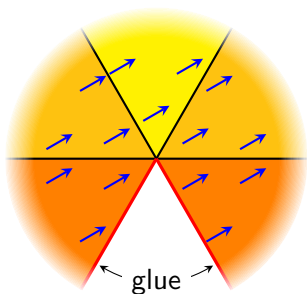
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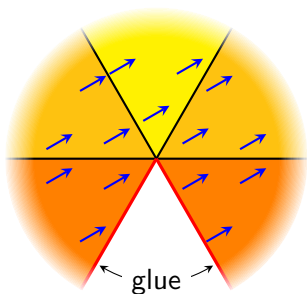


continuous on each triangle

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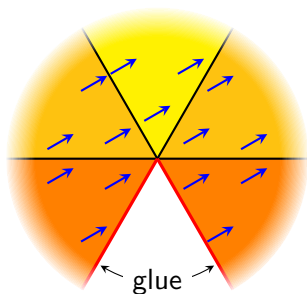


continuous on each triangle
discontinuous across red edge

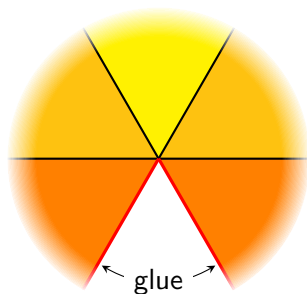
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continuous elements



blow-up elements

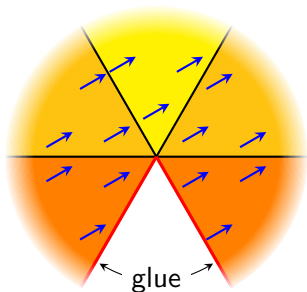


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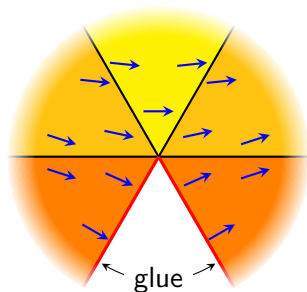
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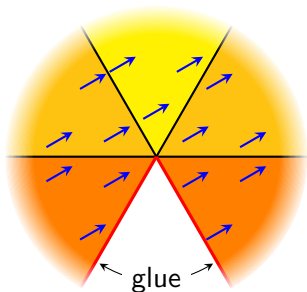


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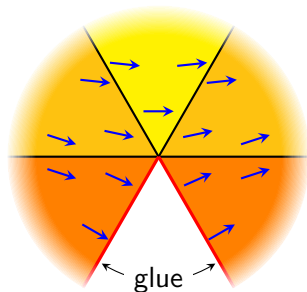
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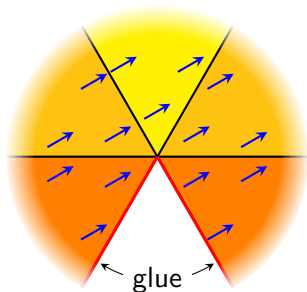


continuous across all edges

Angle defect obstruction to continuous elements

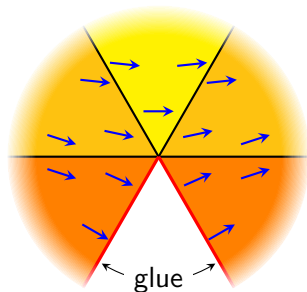
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Vector Laplacian eigenvalue problems

Hodge Laplacian

$$(dd^* + d^*d)v^b = \lambda v^b.$$

- Tangential continuity suffices.
- Standard FEEC works.

Bochner Laplacian

$$\nabla^* \nabla v = \lambda v.$$

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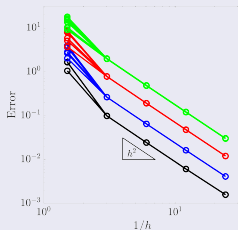
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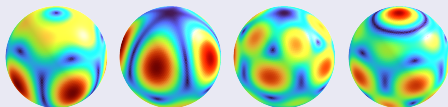
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Bochner Laplacian on sphere using blow-up elements



Eigenvalue error



Eigenfield magnitude
($\lambda = 11, 11, 19, 19$)

There's more

This talk so far

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 - Three radiation sources with rates λ_0 , λ_1 , and λ_2 , sum 1.

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 - Three radiation sources with rates λ_0 , λ_1 , and λ_2 , sum 1.
 - Let t_0 , t_1 , t_2 be the times when the respective radiation sources produce their first particle.

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- A surprising connection to arrival times of Poisson processes, yielding simpler computations.
 - Three radiation sources with rates λ_0 , λ_1 , and λ_2 , sum 1.
 - Let t_0 , t_1 , t_2 be the times when the respective radiation sources produce their first particle.
 - $\frac{\lambda_0 \lambda_1}{\lambda_1 + \lambda_2}$ is the probability that $t_0 \leq t_1 \leq t_2$.

There's more

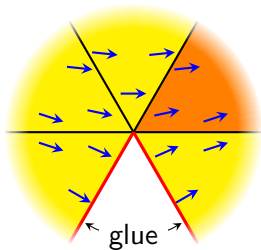
This talk so far

- Lowest order blow-up elements in two dimensions, $b\mathcal{P}_1(T^2)$,
 - including vector fields with components in $b\mathcal{P}_1(T^2)$.

Our preprint

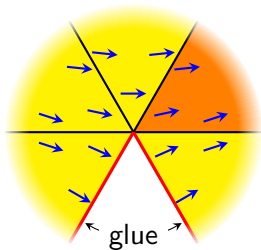
- Differential complex of blow-up Whitney forms, $b\mathcal{P}_1^- \Lambda^k(T^n)$.
 - Shape functions previously studied in (Brasselet, Goresky, MacPherson, 1991), called shadow forms.
- Higher-order blow-up scalar fields $b\mathcal{P}_r(T^n)$.
- A surprising connection to arrival times of Poisson processes, yielding simpler computations.
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- Degrees of freedom in terms of blow-up simplex.

Blowing up



- Even on an individual triangle, the vector field is not continuous at the origin.
- But it is “continuous in polar coordinates,” i.e. in r and θ .

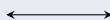
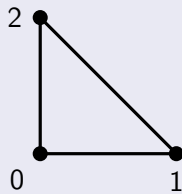
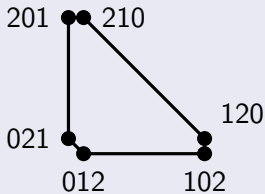
Blowing up



- Even on an individual triangle, the vector field is not continuous at the origin.
- But it is “continuous in polar coordinates,” i.e. in r and θ .




Blowing up manifolds with corners (Melrose, 1996)

- formalizes continuity/smoothness “in polar coordinates”



Smooth “in polar coordinates”

Thank you

-  Yakov Berchenko-Kogan and Evan S. Gawlik
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-  J. P. Brasselet, M. Goresky, and R. MacPherson.
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Amer. J. Math., 113(6):1019–1052, 1991.
-  R. B. Melrose.
Differential analysis on manifolds with corners.
<https://math.mit.edu/~rbm/book.html>, 1996.